

Simulations of Bunch Merging in a Beta Beam Decay Ring

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Abstract. To further study neutrino oscillation properties a Beta Beam facility has been proposed. Beta decaying ions with high kinetic energy are stored in a storage ring ("Decay Ring") with straight sections to create pure focused (anti) electron neutrino beams. However to reach high sensitivity to neutrino oscillation parameters in the experiment the bunched beam intensity and duty cycle in the DR have to be optimized. The first CERN-based scenario, using ${}^6\text{He}$ and ${}^{18}\text{Ne}$ as neutrino sources, has been studied using a bunch merging RF scheme. Two RF cavities at different frequencies are used to capture newly injected bunches and then merge them into the stored bunches. It was shown that this scheme could satisfy the requirements on intensity and duty cycle set by the experiment. This merging scheme has now been revised with new simulation software providing new results for ${}^6\text{He}$ and ${}^{18}\text{Ne}$. Furthermore bunch merging has been studied for the second CERN-based scenario using ${}^8\text{Li}$ and ${}^8\text{B}$.

Keywords: Beta Beams, Decay Ring, Merging, EUROnu-WP4-012

INTRODUCTION

Improving the measurements of the known neutrino oscillation parameters (like Δm_{21}^2 , $|\Delta m_{32}^2|$, θ_{12} and θ_{23}) and determination of the unknown ones (like θ_{13} , $\text{sign}(\Delta m_{32}^2)$ and δ_{cp}) requires precision measurements on a highly pure and intense neutrino beam whose characteristics are well known. One of the proposed next generation neutrino oscillation facilities [1] is the Beta Beam concept [2]. Beta decaying ions are stored at $\gamma = 100$ in a horse-racetrack shaped storage ring, the "Decay Ring". One of the straight sections is aimed at an oscillation experiment and the decaying ions create a highly pure (anti) electron neutrino beam with an opening angle of $1/\gamma$.

Within the EURISOL design study (FP6 [3]) the feasibility of a scenario using ${}^6\text{He}$ (as antineutrino source) and ${}^{18}\text{Ne}$ (as neutrino source) isotopes has been investigated. A 440kt Čerenkov detector located in the Fréjus tunnel at $L_{FP6} = 130$ km distance from CERN was foreseen to detect the incoming neutrinos. The proposed method for bunch injection in the Decay Ring (DR), the so called "RF bunch merging scheme", creates bunches with sufficient intensity while keeping the duty cycle (which coincides with the suppression factor of the experiment) of the DR at 0.58% and well below the sensitivity threshold of 1% given within FP6.

The EUROnu design study (FP7 [4]) includes an additional scenario using ${}^8\text{B}$ (as neutrino source) and ${}^8\text{Li}$ (as antineutrino source) isotopes. Since these isotopes feature a higher Q value, the neutrino energies will be higher and a longer baseline between production and detection is needed. The detector can be placed in Gran Sasso ($L_{FP7} = 732$ km) or Canfranc ($L_{FP7} = 630$ km). A new RF scheme called "barrier buckets" was proposed

but simulations showed that an optimization between bunched intensity and duty cycle could not be achieved [5]. This report therefore focuses on the adaptation of the RF bunch merging scheme used in the EURISOL design study to the isotopes used within EUROnu

Both baselines assume the Beta Beam accelerator complex to be located at CERN, reusing the existing PS and SPS, and that the DR has the same circumference as SPS, namely 6911 m.

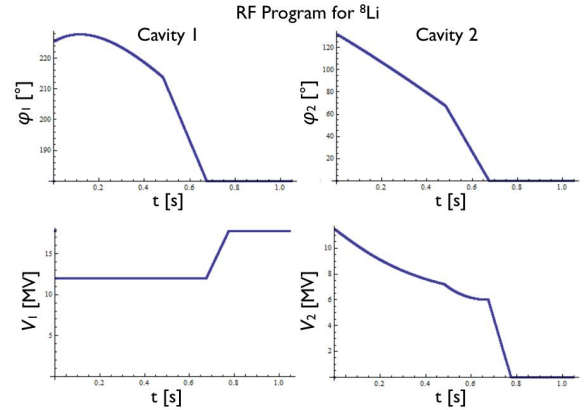


FIGURE 1. Lithium RF program in terms of phases and voltages for both RF systems as used in BBPhase.

RF SIMULATIONS

The RF simulations are done by using the 2D longitudinal phase space program "BBPhase" [6] written to investigate the possibility to use Barrier Buckets in the

Decay Ring [5]. It has since been adapted to perform Bunch Merging simulations. The program tracks the particles on a turn-by-turn basis. After each turn the particles' phases and momenta are evaluated dependent on the voltage they have seen from the RF system in the previous turn. The combined voltage from two cavities with frequencies $f_1 = 40$ MHz ($h_1 = 924$) and $f_2 = 80$ MHz ($h_2 = 2h_1 = 1848$), is

$$V_{rf} = \sum_{i=1}^2 V_i \sin(h_i \varphi + h_i \varphi_i) \quad (1)$$

where the (time dependent) maximum voltages are V_i , the RF phase when a reference particle passes through the long wavelength RF cavity is $-h_1 \varphi_1$ and $-h_2 \varphi_2$ when it passes through the fast frequency RF cavity and φ is the azimuthal difference between the reference particle and a given particle in the bunch.

The merging program was theoretically optimized for ${}^6\text{He}$ and ${}^{18}\text{Ne}$ [7] and has now been adapted for ${}^8\text{Li}$ and ${}^8\text{B}$. The plots in fig. 1 represent the changes in phases and voltages for the merging program of Lithium.

We apply the merging program (fig. 1) of Lithium to the case that one bunch is already injected and circulating in the main bucket (black dots in fig. 2 represent multiparticles of this bunch). The newly injected beam (red) is distributed in a longitudinal phase space ellipse at a momentum excess of 5 per mil (fig. 2a). The best capture efficiency could be obtained by slightly lowering the particles momentum excess to 4.92 per mil and moving it about 9.6 cm behind the stored bunch. The particles perform a quarter synchrotron turn and loose their elliptical distribution since the movement around the synchronous particle is not linear in this region of phase space (fig. 2b). This causes some particles to be outside the elliptical shaped capture bucket when the second cavity is turned on. Those particles are not captured and circle on the separatrix around both buckets. Asymmetric merging is applied during 0.5 seconds where the main bucket shrinks in size (until both buckets have the same size, fig. 2c, 2d and 2e). After that the symmetric merging procedure takes about 0.2 seconds to decrease the distance between the two buckets to zero (fig. 2f and 2g). At this point the second cavity is turned off and the main RF is progressively tuned to nominal voltage. Before the next injection the particles circulating around the bucket have to be collimated at 2.5 per mil momentum excess (fig. 2h) or they will hit the septum blade during the next injection.

These simulations have been done with all FP6 and FP7 ions together with capture efficiency studies. Capture efficiency is defined as quotient between not-collimated (e.g. running or decayed) particles and injected particles at the end of each SPS cycle. It was possible to reach an efficiency of slightly below 90% in all cases (table 1). However the results for ${}^8\text{Li}$ and ${}^8\text{B}$

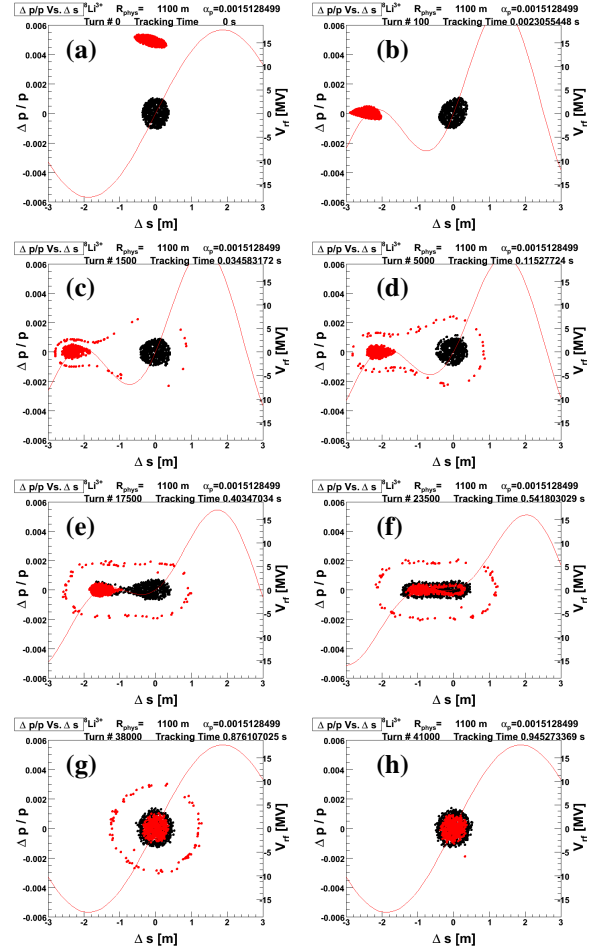


FIGURE 2. Different phases of the RF Bunch Merging procedure. The incoming bunch is injected (a) and then captured by the switched-on second cavity after a quarter turn (b). Asymmetric (c, d, e) and symmetric (f, g) merging are applied and the beam is collimated before the end of the procedure (h).

are very preliminary since no data on longitudinal emittances was available at the time of the simulations.

During the merging process the longitudinal emittance of the stored beam is increased due to RF gymnastics. The collimation at 2.5 per mil momentum excess will therefore limit the bunch size after a certain number of injections. Together with radioactive decay (also included in the BBPhase simulation) the collimation will therefore cause the accumulated number of ions in the bunch to saturate. We obtain the number of particles after which a bunch is saturated by repeatedly applying the merging program with the appropriate cycle times for each ion (table 1). The amount of remaining particles per bunch is plotted every 4000 turns over a period of 25 injections (27 for ${}^{18}\text{Ne}$ due to later saturation of Neon) in figure 3. Due to uncertainties in the ion production (R&D is ongoing) the amount of injected ions per cycle is an es-

TABLE 1. DR parameters and simulation results for all isotopes. Nominal ν -rates for ^8Li and ^8B are 5 times bigger than for ^6He and ^{18}Ne respectively [8]. Results for Lithium and Boron are preliminary. (Acc. stands for accumulated.)

	^6He	^{18}Ne	^8Li	^8B
$t_{1/2}$ at rest [ms]	807	1872	838	770
SPS cycle time [s]	6.0	3.6	4.8	3.6
Source rate [$10^{13}/\text{s}$]	2	2.09	9	9
Inject./Bunch [10^{11}]	4.87	2.35	21.50	8.43
Capture Efficiency	88.8%	87.8%	88.2%	89.0%
Acc./Bunch [10^{11}]	40	31	173	74
Acc. in DR [10^{13}]	8	6.20	35	13.6
ν -rate [$10^{18}/\text{year}$]	2.4	0.92	10.2	4.34
Nom. ν -rate [10^{18}]	2.9	1.1	14.5	5.5
ν -rate ratio	0.828	0.836	0.705	0.788

timated value for all isotopes.

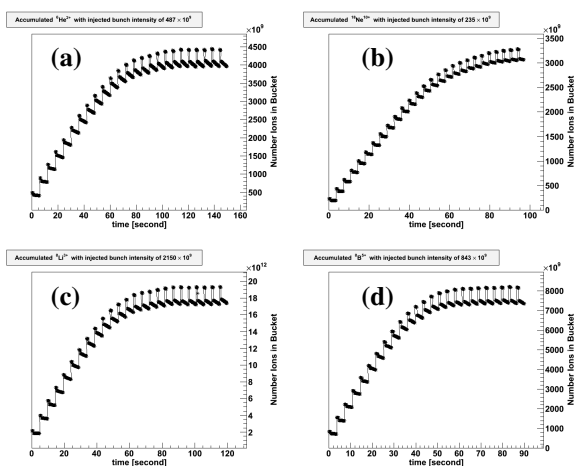


FIGURE 3. Accumulation of (a) ^6He , (b) ^{18}Ne , (c) ^8Li and (d) ^8B ions per bunch.

At saturation the size of one stored bunch is about 2 meters for all ions. The duty cycle of the DR (= suppression factor of the experiment), assuming 20 injected bunches, is then $DC = 20 \cdot 2m/6911m = 0.58\%$. This holds for all tested ions which is in agreement with the results in [7] and extends them to ^8Li and ^8B . From the number of accumulated ions and 20 injected bunches we can calculate the annual neutrino flux via a Mathematica notebook [9]. It turns out that the amount of ions stored at saturation ($3.1 \cdot 10^{12}\text{Ne}$ per bunch and $4 \cdot 10^{12}\text{He}$ per bunch) corresponds to only 84% (83%) of the nominal (anti) neutrino flux [10] (see table 1). In the case of ^8Li and ^8B sensitivity plots shows that nominal neutrino rate needed is 5 times larger than for ^6He and ^{18}Ne [8]. The (anti) neutrino flux achieved from the number of ions stored in the DR at saturation corresponds to 71% (^8Li) and 79% (^8B) of the nominal fluxes.

We are however well below the requirements of the duty cycle which is at maximum 1%.

CONCLUSIONS

It was possible to adapt the radio frequency bunch merging scheme to the 2D program BBPhase and recreate the capture efficiencies already calculated in case of FP6 ions [7]. These results could also be extended to ^8Li and ^8B using preliminary values for emittances and ion production showing the feasibility of bunch merging for these isotopes. Moreover the duty factor considering 20 accumulated bunches in the Decay Ring was found to be 0.58% which also agrees with the previous theoretical results and is well below the upper limit of 1%. In light of these results the Bunch Merging procedure is at present the preferred method to create ion bunches with a sufficiently low duty cycle inside the Decay Ring.

However the amount of ions that could be stored in one bunch at saturation was found to be lower than needed for the production of nominal (anti) neutrino flux in the Helium and Neon case. For the Lithium and Boron case further R&D is needed in terms of ion production and transportation. The simulations will then have to be repeated with new values for longitudinal emittances and injected bunch intensities.

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