

# NuFact10

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## Theoretical highlights on neutrino-nucleus interactions: current challenges

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# Introduction

- Neutrino-nucleus interactions are **important** for:
  - Oscillation experiments
    - $\nu$  detection,  $E_\nu$  reconstruction,  $\nu$  flux calibration
    - Electron-like backgrounds:
      - NC  $\pi^0$  production (incoherent, coherent)
      - Photon emission in NC
  - Hadronic physics
    - Nucleon and Nucleon-Resonance ( $N-\Delta$ ,  $N-N^*$ ) axial form factors
    - **Strangeness** content of the nucleon spin
  - Nuclear physics
    - Information about: nuclear correlations, MEC, spectral functions
    - **nuclear effects**: essential for the interpretation of the data

# Introduction

- Most relevant processes in the few-GeV region:
  - Quasielastic scattering
  - Single pion production (incoherent and coherent)

# QE scattering

- $\nu$  detection:  $\nu_\mu n \rightarrow \mu^- p$
- $E_\nu$  reconstruction:
  - Assumes that  $E_\nu = \frac{2m_n E_\mu - m_\mu^2 - m_n^2 + m_p^2}{2(m_n - E_\mu + p_\mu \cos \theta_\mu)}$
  - Important for oscillations:  $P(\nu_\mu \rightarrow \nu_\tau) = \sin^2 2\theta_{23} \sin^2 \frac{\Delta m_{23}^2 L}{2E_\nu}$
- QE  $\sigma$  affect the expected sensitivities to oscillation parameters
  - Example: Fernandez, Meloni, arXiv:1010.2329
  - $\beta$  beam hypothetical exp.:  $\langle E_\nu \rangle \sim 0.3 \text{ GeV}$ ,  $^{16}\text{O}$  target
  - Sensitivities to  $\theta_{13}$  and  $\delta_{CP}$  vary  $\sim 10\text{-}30\%$  depending on the nuclear CCQE model

# $\nu$ QE scattering

- The (CC) elementary process:  $\nu_\mu(k)$   $n(p) \rightarrow \mu^-(k')$   $p(p')$

$$\mathcal{M} = \frac{G_F \cos \theta_C}{\sqrt{2}} l^\alpha J_\alpha$$

where  $l^\alpha = \bar{u}(k') \gamma^\alpha (1 - \gamma_5) u(k)$

$$J_\alpha = \bar{u}(p') \left[ \gamma_\alpha F_1^V + \frac{i}{2M} \sigma_{\alpha\beta} q^\beta F_2^V + \gamma_\mu \gamma_5 F_A + \frac{q_\mu}{M} \gamma_5 F_P \right] u(p)$$

- Vector form factors:  $F_{12}^V = F_{12}^p - F_{12}^n$  extracted from e-p, e-d data

- Axial form factors:

$$F_A(Q^2) = g_A \left( 1 + \frac{Q^2}{M_A^2} \right)^{-2} \quad \text{dipole ansatz}, \quad F_P(Q^2) = \frac{2M^2}{Q^2 + m_\pi^2} F_A(Q^2) \quad \text{PCAC}$$

# $\nu$ QE scattering

- The (CC) elementary process:  $\nu_\mu(k)$   $n(p) \rightarrow \mu^-(k')$   $p(p')$

$$F_A(Q^2) = g_A \left( 1 + \frac{Q^2}{M_A^2} \right)^{-2}$$

- $g_A = 1.267 \leftarrow \beta$  decay
- $M_A = 1.016 \pm 0.026$  GeV ( $\nu d$ ,  $\bar{\nu} p$ ) Bodek et al., EPJC 53 (2008)
- $M_A$  from  $\pi$  electroproduction on p:
  - Connected to  $F_A$  at threshold and in the chiral limit ( $m_\pi = 0$ )
  - Using models to connect with data  $\Rightarrow$
  - $M_A^{ep} = 1.069 \pm 0.016$  GeV Liesenfeld et al., PLB 468 (1999) 20
  - A more careful evaluation in ChPT Bernard et al., PRL 69 (1992) 1877

$$\langle r_A^2 \rangle_e = \langle r_A^2 \rangle_\nu + \frac{3}{64f_\pi} \left( 1 - \frac{12}{\pi^2} \right), \langle r_A^2 \rangle = \frac{12}{M_A^2}$$

- $M_A = M_A^{ep} - \Delta M_A$ ,  $\Delta M_A = 0.055$  GeV  $\Rightarrow M_A = 1.014$  GeV

# $\nu$ QE scattering

- Relativistic Global Fermi Gas **Smith, Moniz, NPB 43 (1972) 605**

- Impulse Approximation

- Fermi motion  $f(\vec{r}, \vec{p}) = \Theta(p_F - |\vec{p}|)$

- Pauli blocking  $P_{\text{Pauli}} = 1 - \Theta(p_F - |\vec{p}|)$

- Average binding energy  $E = \sqrt{\vec{p}^2 + m_N^2} - \epsilon_B$

- Explains the main features of the  $(e, e')$  **inclusive**  $\sigma$  in the **QE** region

- **Fails** in the details (nuclear dynamics needed)

# $\nu$ QE scattering

- Spectral functions of nucleons in nuclei
  - The nucleon propagator can be cast as

Benhar et al., PRD 72 (2005)  
Ankowski, Sobczyk, PRC 67 (2008)  
Nieves et al., PRC 70 (2004)  
Leitner et al., PRC 79 (2009)

$$G(p) = \int d\omega \frac{S_h(\omega, \vec{p})}{p^0 - \omega - i\epsilon} + \int d\omega \frac{S_p(\omega, \vec{p})}{p^0 - \omega + i\epsilon}$$

- $S_{h(p)}$  ← hole (particle) spectral functions: 4-momentum ( $p$ ) distributions of the struck (outgoing) nucleons

$$S_{p,h}(p) = -\frac{1}{\pi} \frac{\text{Im}\Sigma(p)}{[p^2 - M^2 - \text{Re}\Sigma(p)]^2 + [\text{Im}\Sigma(p)]^2}$$

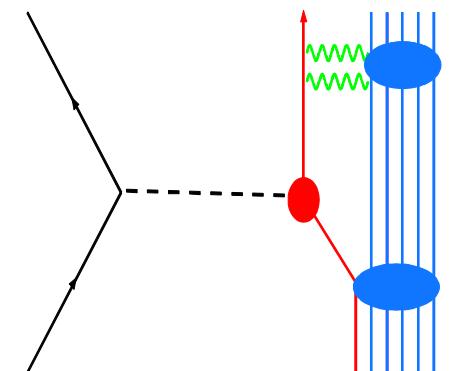
- $\Sigma$  ← nucleon selfenergy
- Better description of (e,e') inclusive  $\sigma$

# $\nu$ QE scattering

- Relativistic mean field
  - Impulse Approximation
  - Initial nucleon in a **bound state** (shell)
    - $\Psi_i$  : Dirac eq. in a mean field potential ( $\omega$ - $\sigma$  model)
  - Final nucleon
    - PWIA
    - RDWIA:  $\Psi_f$  : Dirac eq. for scattering state
    - Glauber
  - **Problem:** nucleon absorption that reduces the c.s.
  - Can be used to study:
    - 1N knockout
    - inclusive processes: with **only**  $\text{Re}[V_{\text{opt}}]$

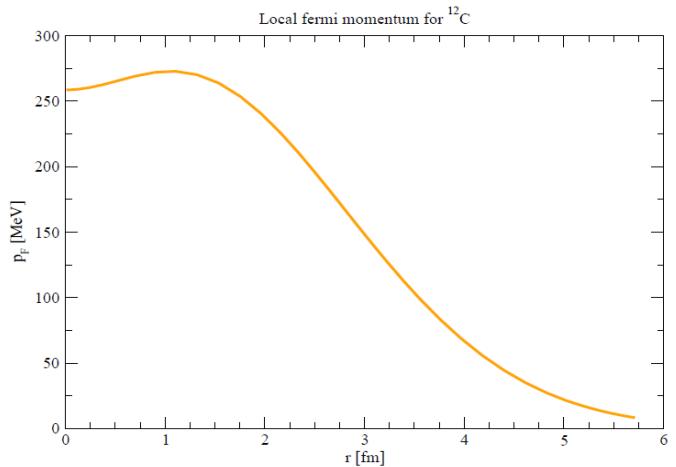
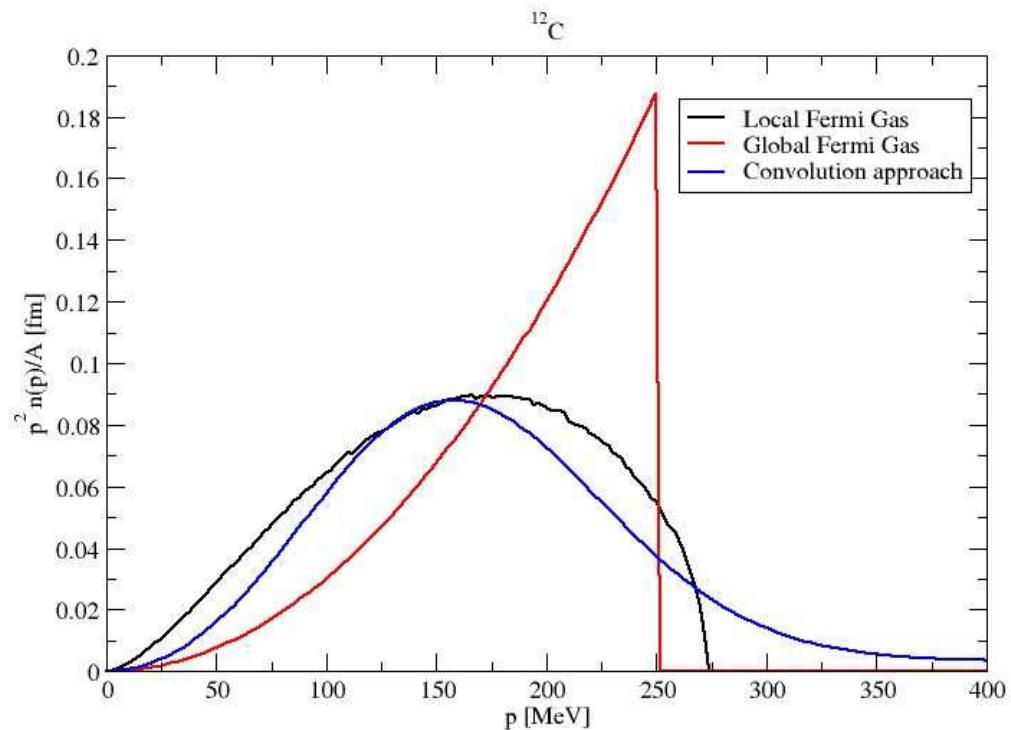
Martinez et al., PRC 73 (2006)  
Budkevich, Kulagin, PRC 76 (2007)

} Complex optical potential



# $\nu$ QE scattering

- Local Fermi Gas  $p_F(r) = [\frac{3}{2}\pi^2\rho(r)]^{1/3}$

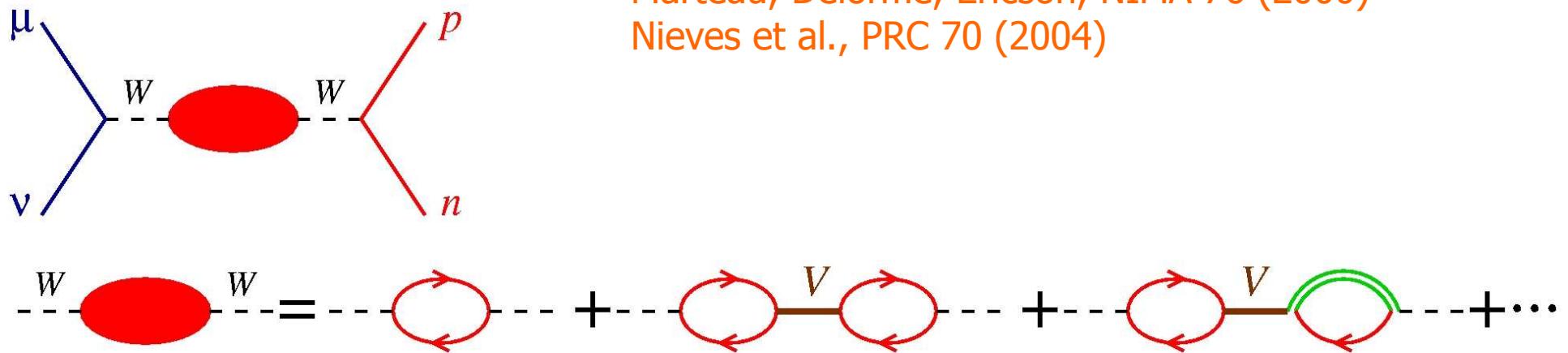


Convolution model:  
Ciofi degli Atti, Simula, PRC 53 (1996)

- Space-momentum correlations **absent** in the GFG
- Reasonable for medium/heavy nuclei
- Microscopic many-body effects are tractable  
(calculations in infinite nuclear matter)

# $\nu$ QE scattering

## ■ RPA long range correlations



Singh, Oset, NPA 542 (1992)

Marteau, Delorme, Ericson, NIMA 76 (2000)

Nieves et al., PRC 70 (2004)

## ■ Incorporates $N$ -hole and $\Delta$ -hole states

## ■ $V$ : $\pi$ , $\rho$ exchange, Landau-Migdal parameter $g'$

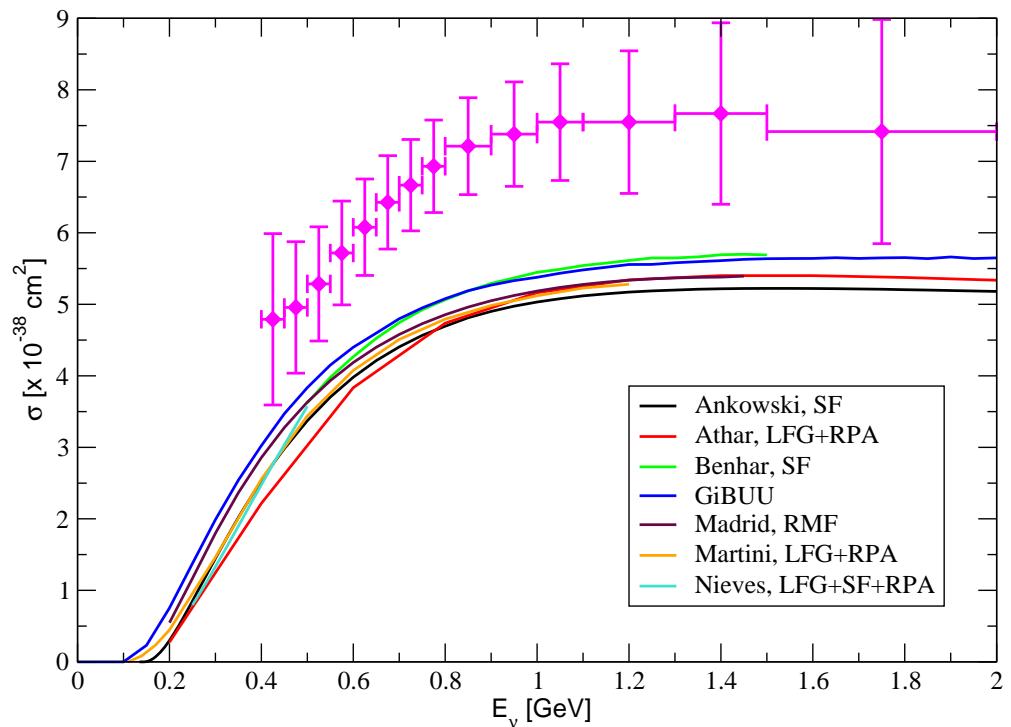
## ■ Describes correctly $\mu$ capture on $^{12}\text{C}$ and LSND CCQE

## ■ **Collective effect:** important at low $Q^2$ for $\nu$ QE

# $\nu$ QE scattering

## ■ Comparison to MiniBooNE $\sigma$

CCQE on  $^{12}\text{C}$



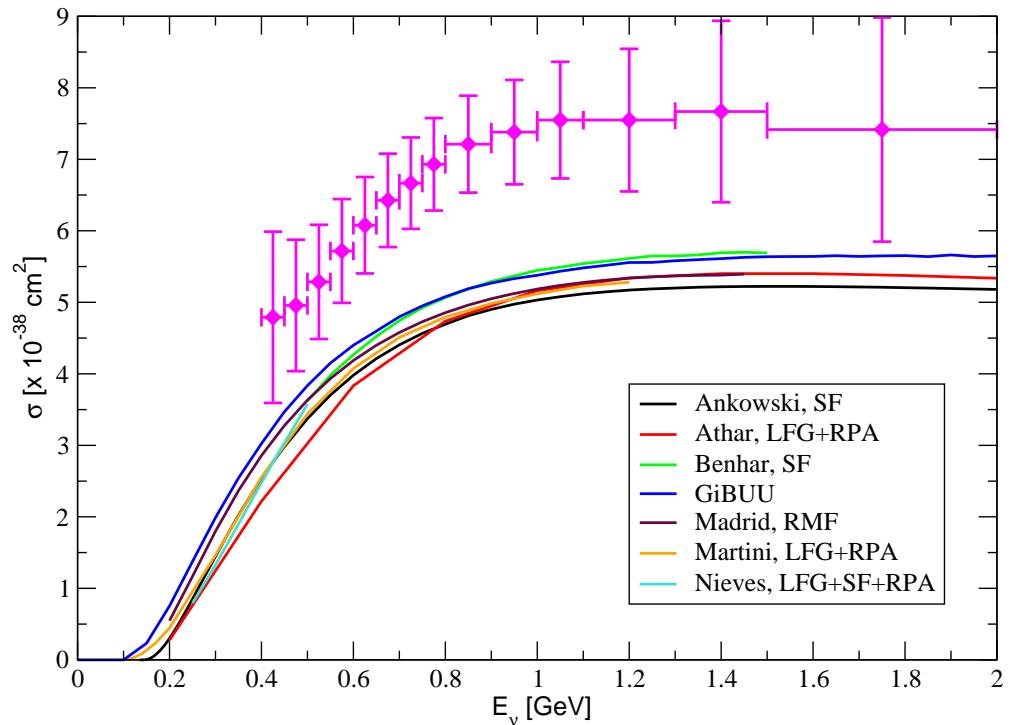
Source: Boyd et al., AIP Conf. Proc. 1189  
Data: MiniBooNE, PRD 81, 092005 (2009)

- At  $E_\nu = 0.8 \text{ GeV}$ :  $\sigma_{th} \sim 4.5-5 < \sigma_{MB} \sim 7 \times 10^{-38} \text{ cm}^2$
- CCQE models with  $M_A \sim 1 \text{ GeV}$  **cannot** reproduce MiniBooNE  $\sigma$

# $\nu$ QE scattering

## ■ Comparison to MiniBooNE $\sigma$

CCQE on  $^{12}\text{C}$

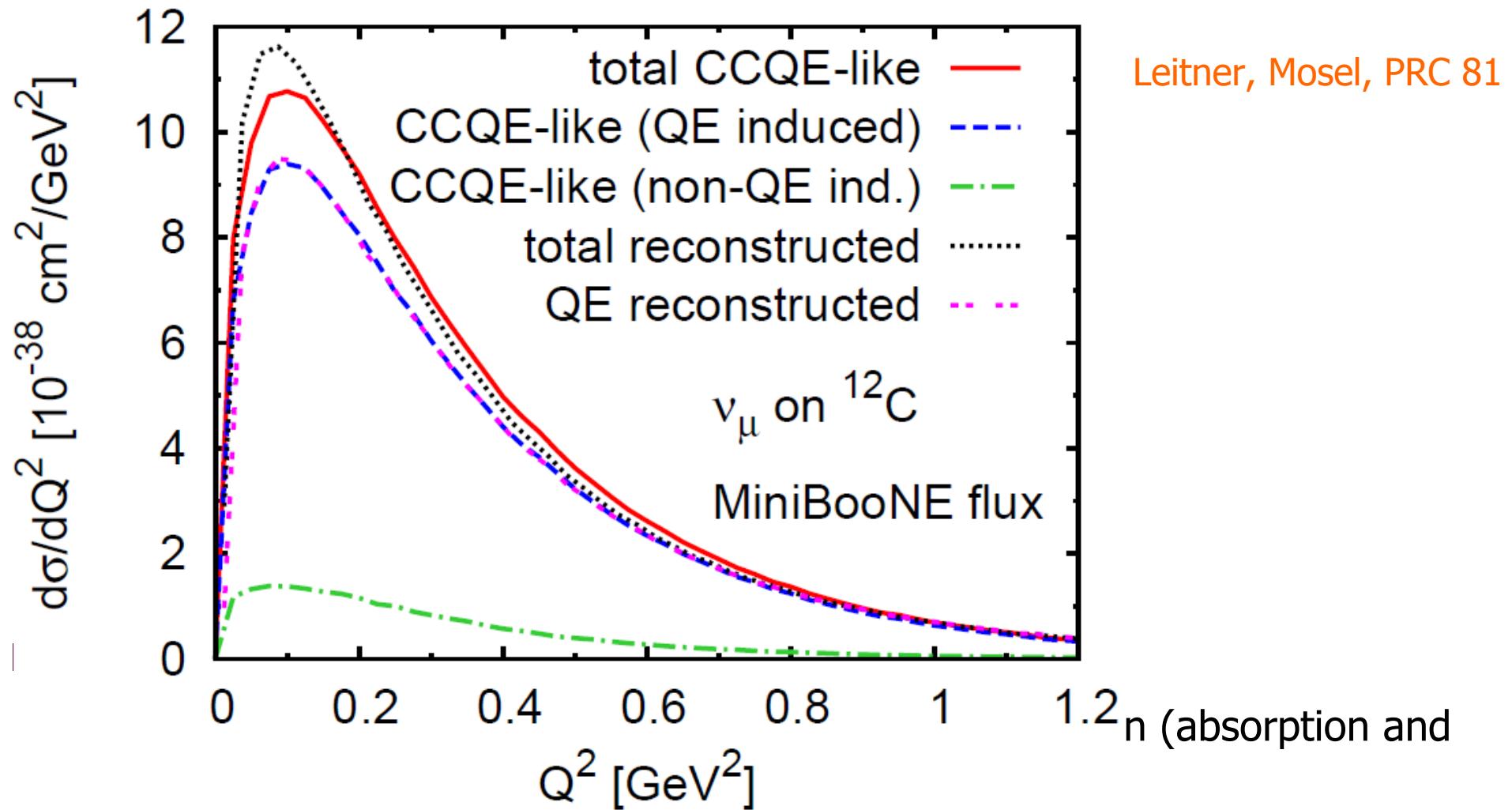


Source: Boyd et al., AIP Conf. Proc. 1189  
Data: MiniBooNE, PRD 81, 092005 (2009)

- Background (CCQE-like) depends on the  $\pi$  propagation (absorption and charge exchange) model (NUANCE)
- $E_\nu$  reconstruction (unfolding)

# $\nu$ QE scattering

- Comparison to MiniBooNE  $\sigma$



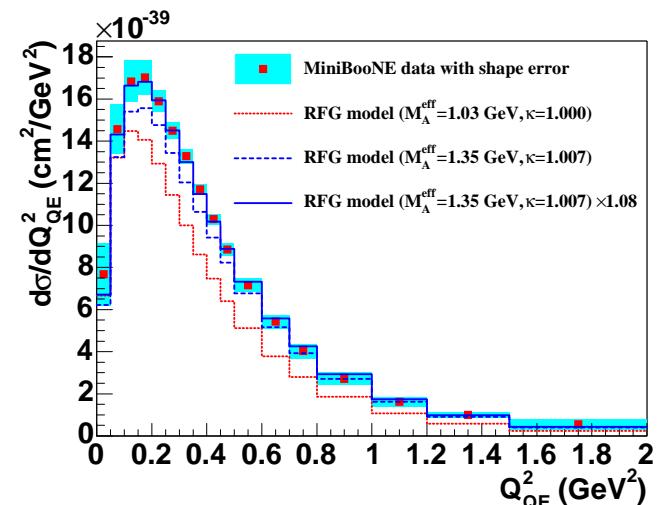
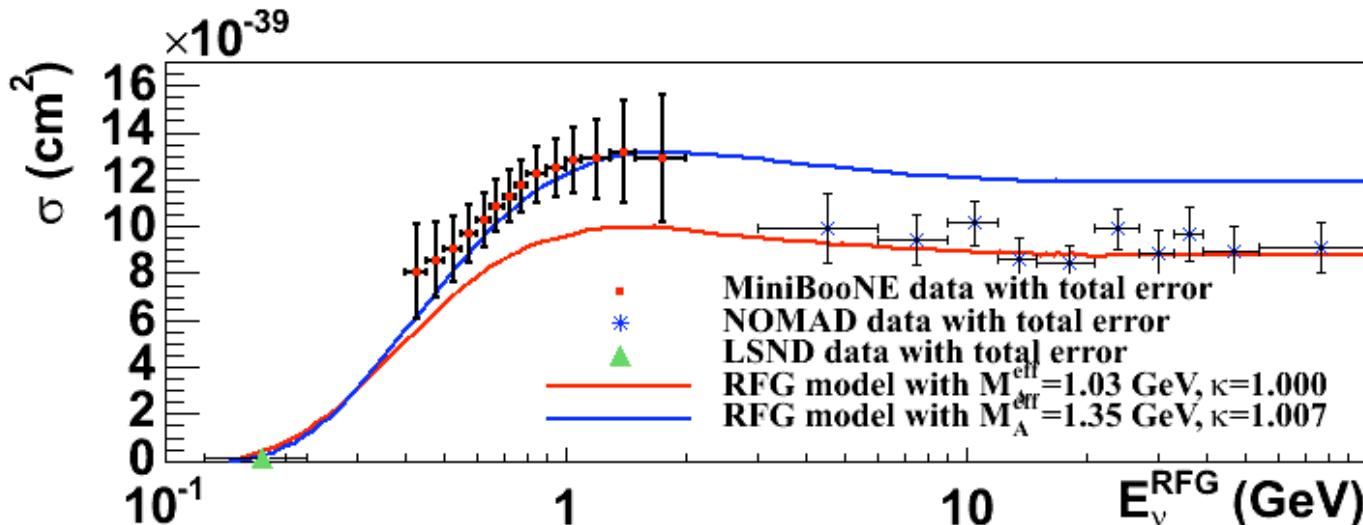
- $E_\nu$  reconstruction (unfolding): good if the CCQE-like background is properly subtracted

# $\nu$ QE scattering

- Comparison to MiniBooNE  $\sigma$

- Proposed solutions:

- $M_A = 1.35 \pm 0.17 \text{ GeV}$  (RFG) [MiniBooNE, PRD 81, 092005 \(2010\)](#)



- However,  $M_A > 1 \text{ GeV}$  is **incompatible** with:

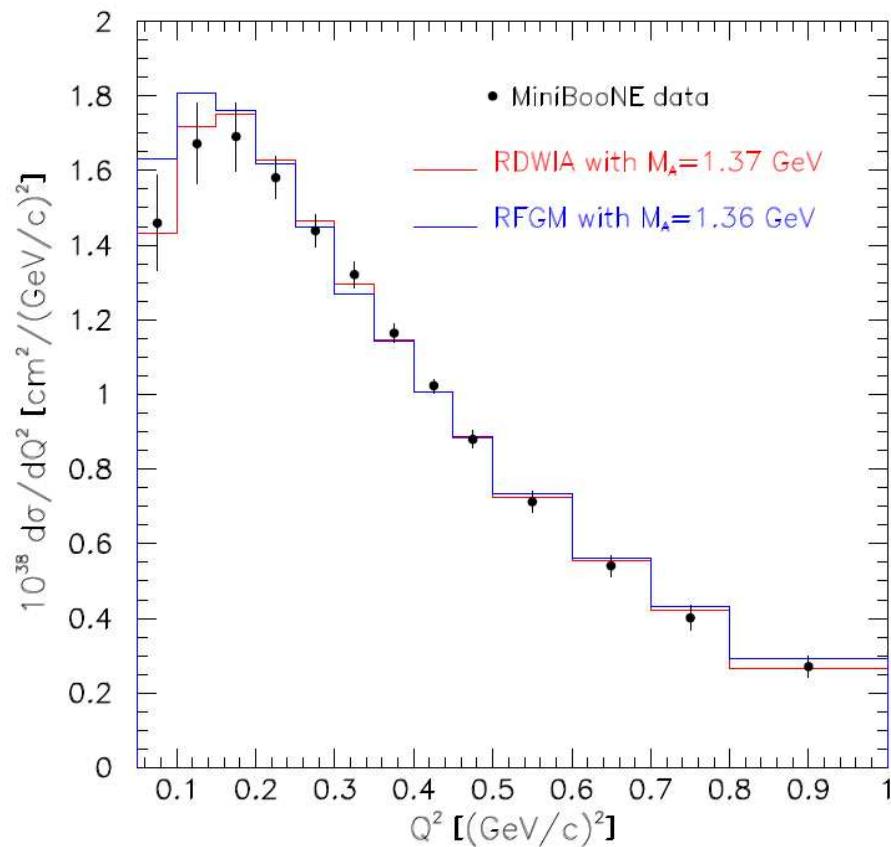
- $\nu d, \bar{\nu} p$  data
    - $\pi$  electroproduction on p (at low  $Q^2$ )
    - NOMAD:  $M_A = 1.05 \pm 0.02(\text{stat}) \pm 0.06(\text{sys}) \text{ GeV}$  [Lyubushkin et al., EPJ C63](#)

# $\nu$ QE scattering

## ■ Comparison to MiniBooNE $\sigma$

### ■ Proposed solutions:

- $M_A = 1.37 \text{ GeV}$  (RDWIA) Butkevich, arXiv:1006.1595
- Fit to  $d\sigma/dQ^2$  (shape only)



- Better than RFG at low  $Q^2$

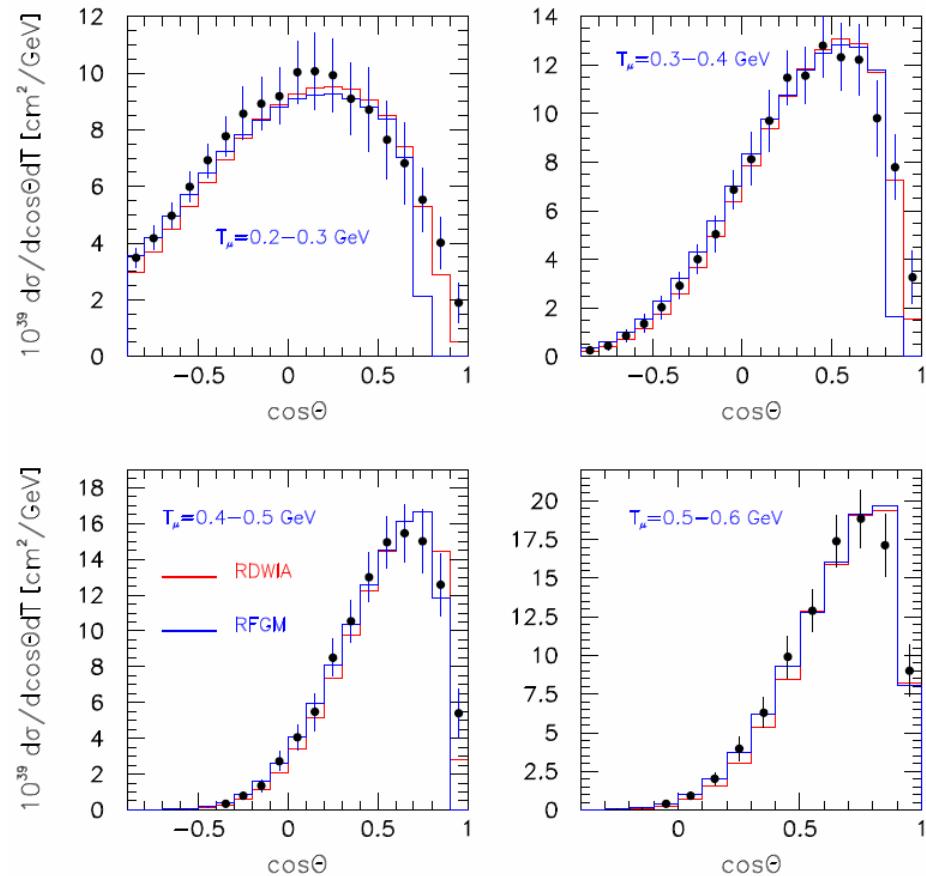
# $\nu$ QE scattering

## ■ Comparison to MiniBooNE $\sigma$

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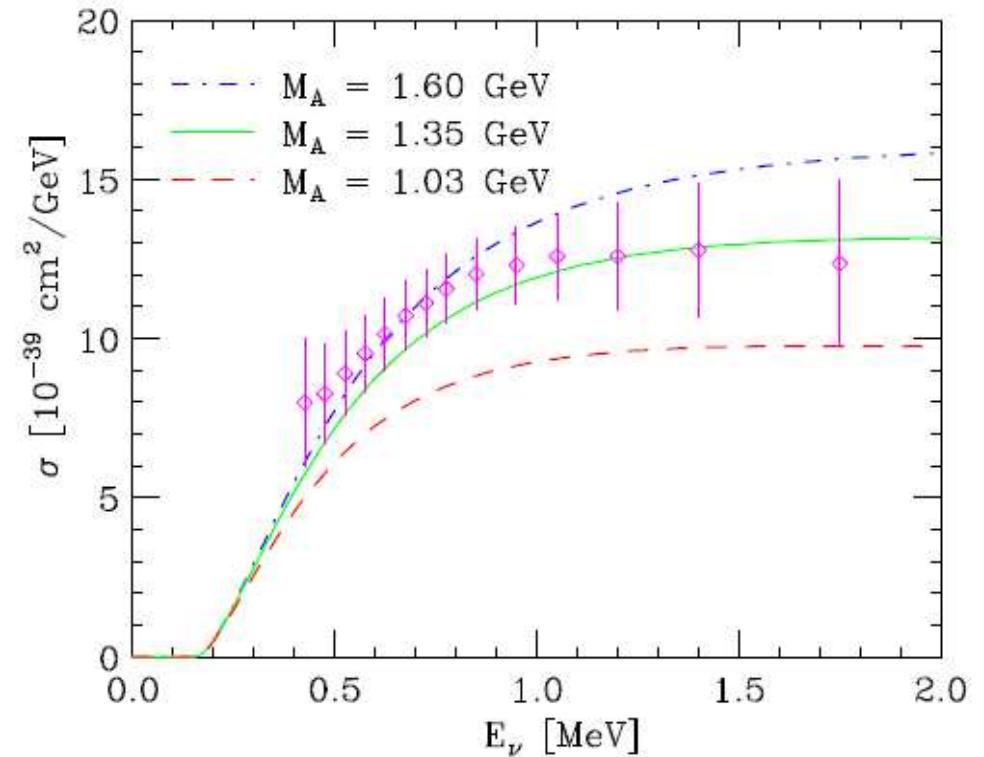
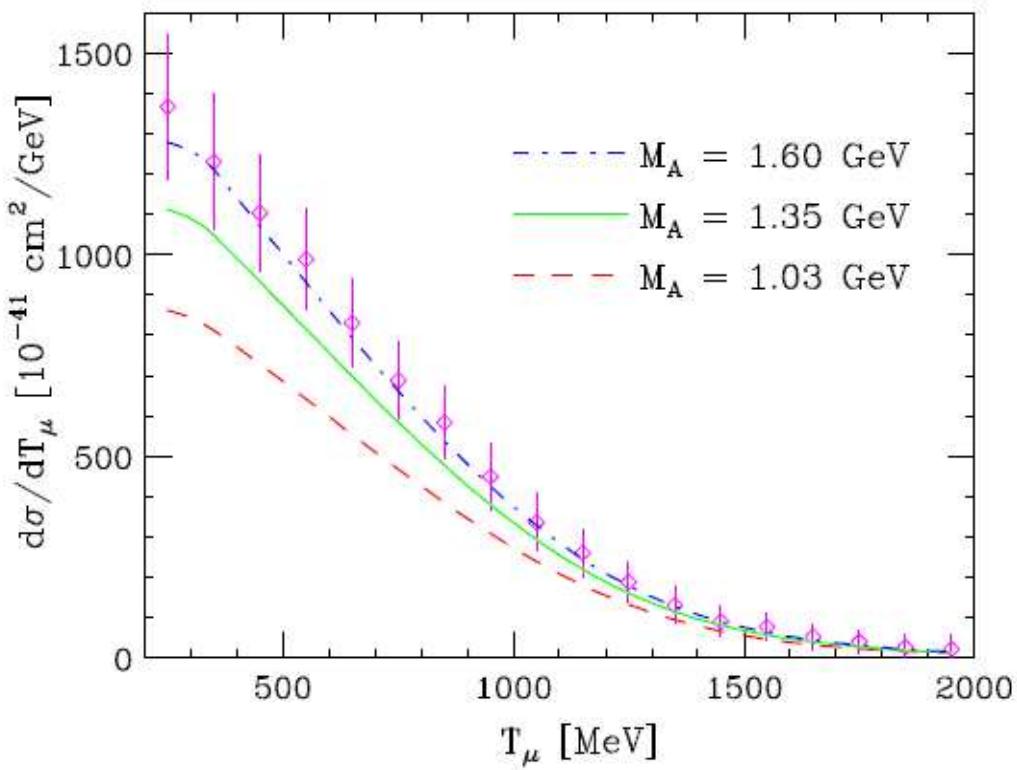
- Good description of double differential cross section

# $\nu$ QE scattering

## ■ Comparison to MiniBooNE $\sigma$

### ■ Proposed solutions:

- $M_A = 1.6$  GeV (Spectral Function) Benhar, Coletti, Meloni, PRL 105 (2010)
- Fit to  $d\sigma/dQ^2$

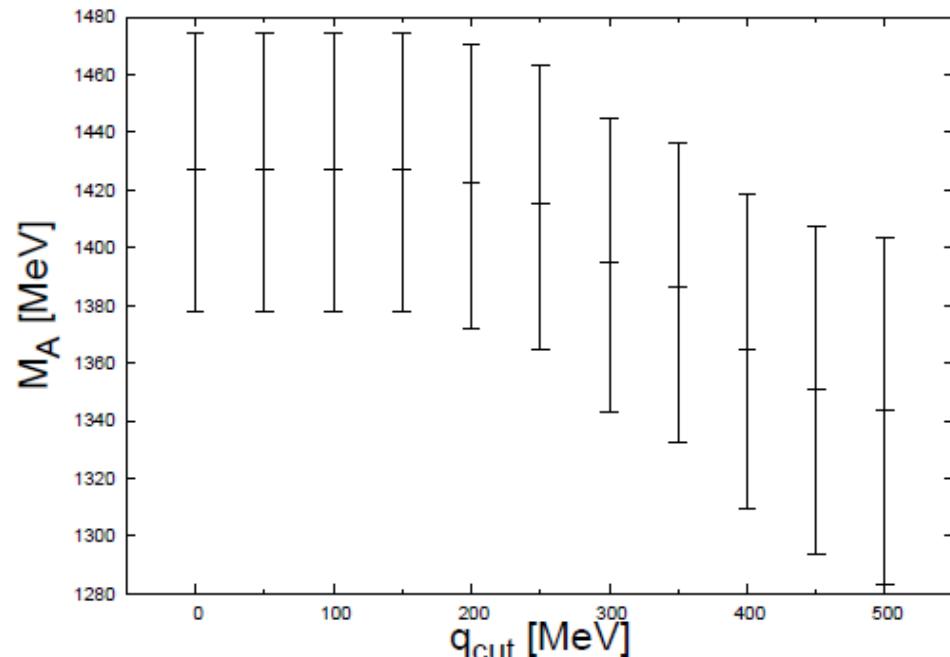


# $\nu$ QE scattering

## ■ Comparison to MiniBooNE $\sigma$

### ■ Proposed solutions:

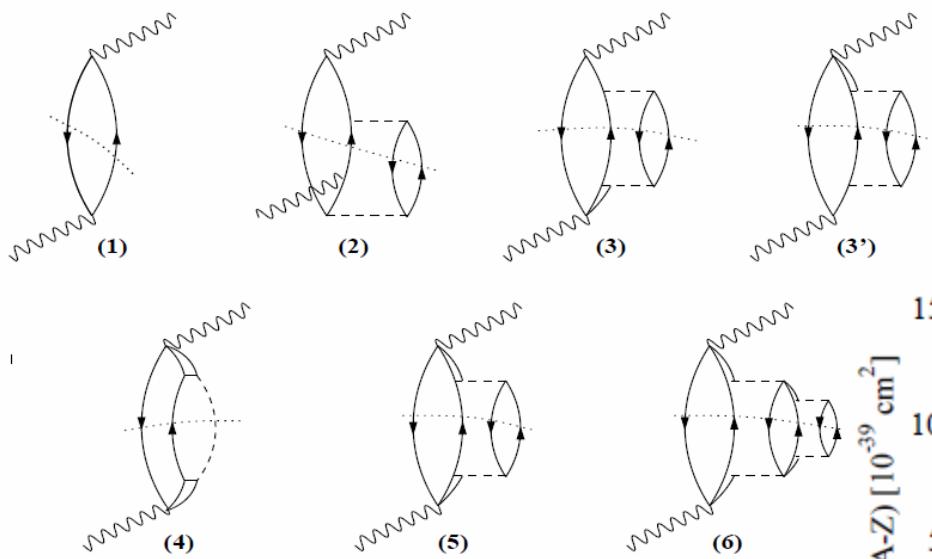
- $M_A = 1.343 \pm 0.060$  GeV (Spectral Function) Juszczak, Sobczyk, Zmuda, arXiv:1007.2195
- Fit to double differential c. s. including flux uncertainty
- Momentum transfer cut  $q_{\text{cut}} = 500$  MeV to exclude IA breakdown region: insufficient to reconcile MiniBooNE with exp. on deuterium



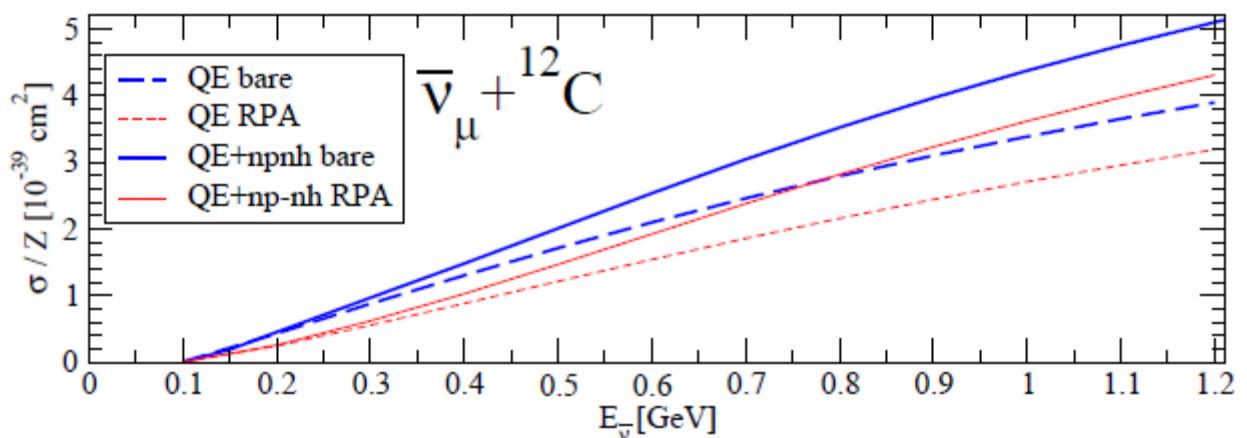
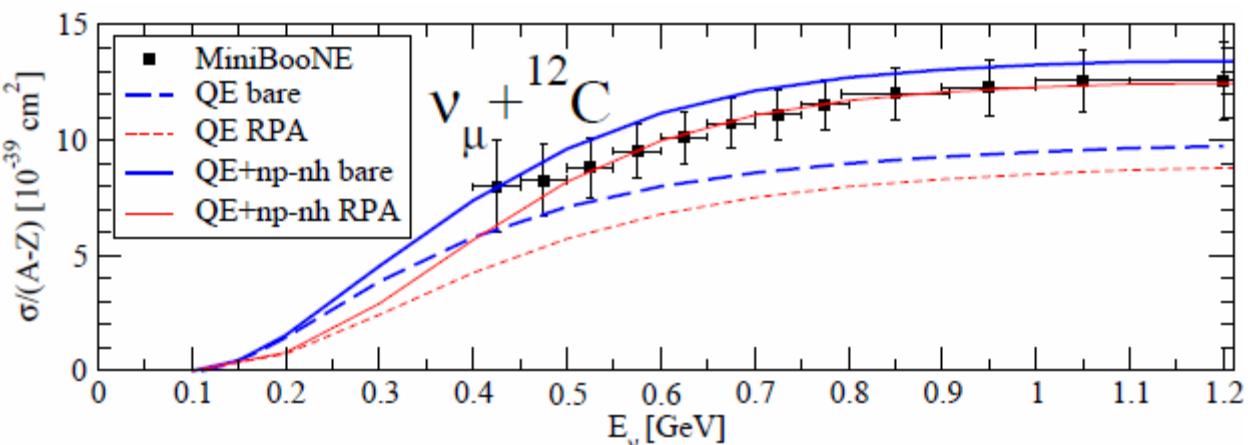
# $\nu$ QE scattering

## ■ Comparison to MiniBooNE $\sigma$

### ■ Many body RPA Martini et al., PRC 80 (2009), 81 (2010)



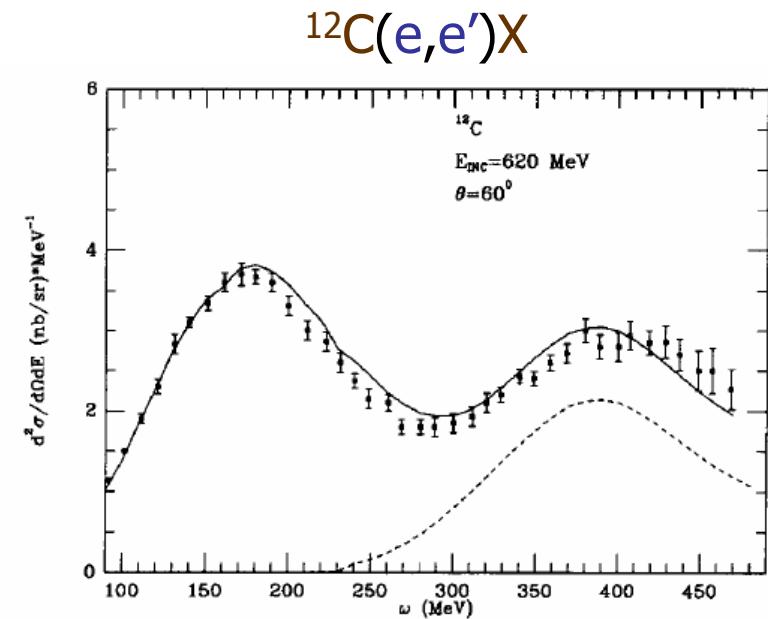
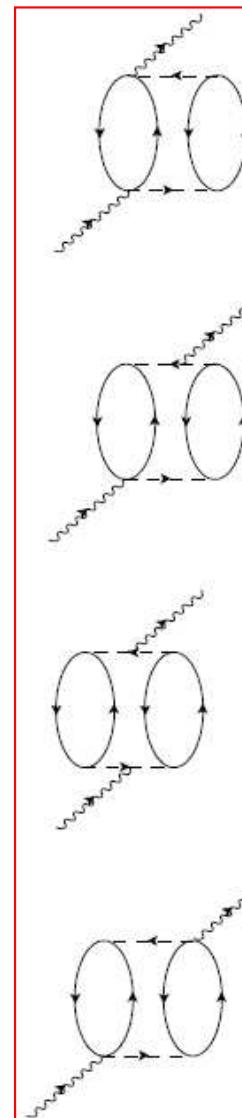
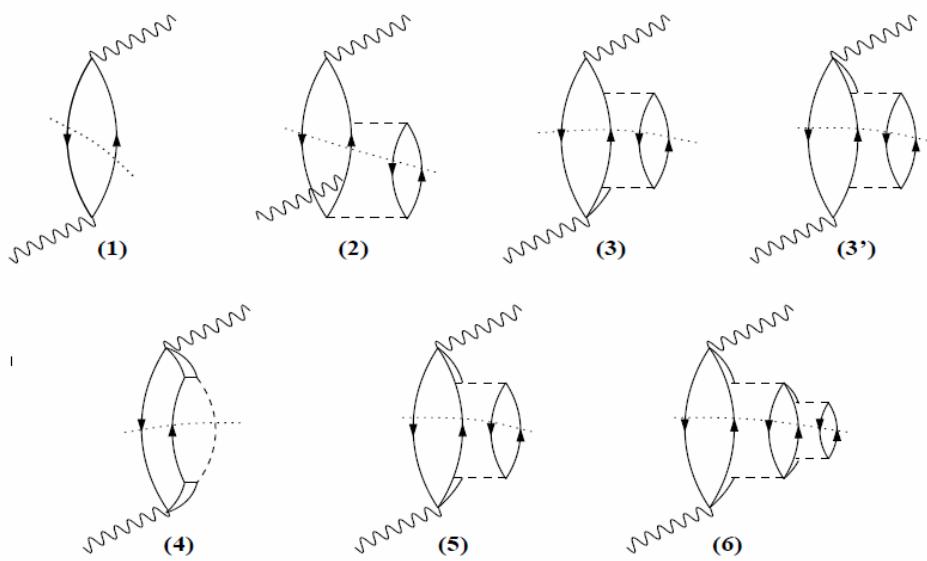
- RPA: small reduction
- Large 2p-2h contribution to  $\nu + ^{12}\text{C}$  mainly from (2),(3),(3')



# $\nu$ QE scattering

## ■ Comparison to MiniBooNE $\sigma$

### ■ Many body RPA Martini et al., PRC 80 (2009), 81 (2010)



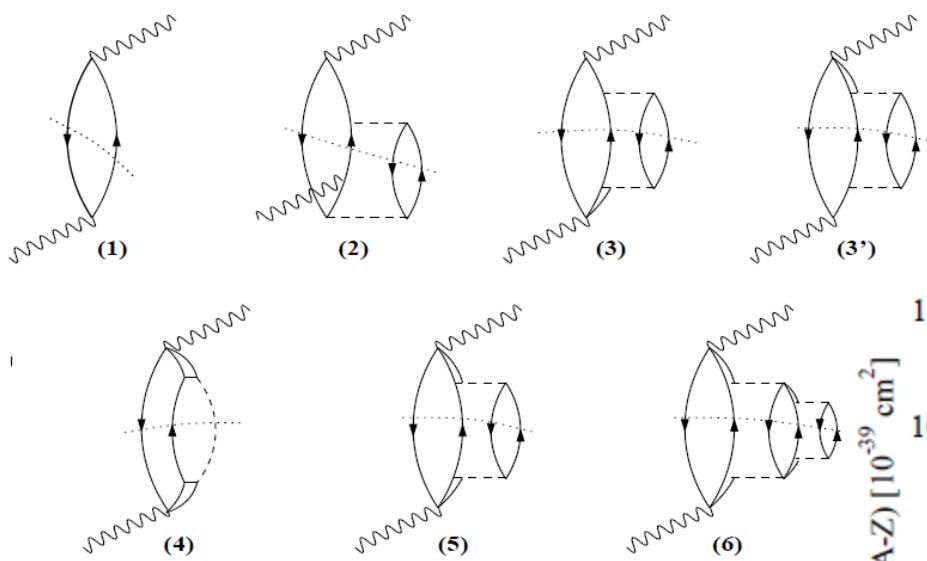
- RPA: small reduction
- Large 2p-2h contribution to  $\nu$   $^{12}\text{C}$  mainly from (2),(3),(3')
- MEC absent (but important for  $(\text{e},\text{e}')$  in the dip region)

Gil, Nieves, Oset, NPA627

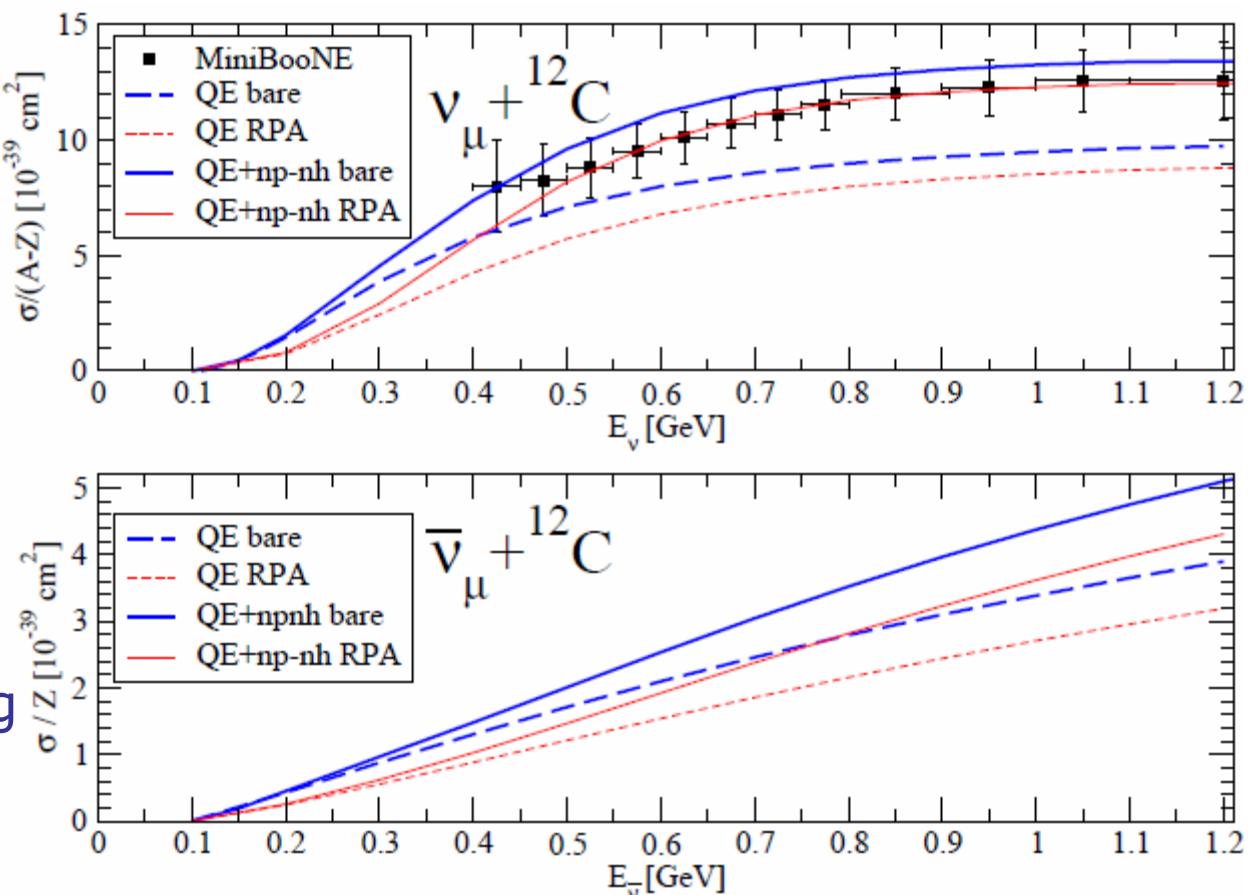
# $\nu$ QE scattering

## ■ Comparison to MiniBooNE $\sigma$

■ Many body RPA Martini et al., PRC 80 (2009), 81 (2010)



- RPA: small reduction
- Large 2p-2h contribution to  $\bar{\nu} + ^{12}\text{C}$  mainly from (2),(3),(3')
- Prediction: smaller 2p-2h contribution to anti- $\bar{\nu} + ^{12}\text{C}$
- Detailed tests against e-scattering data are necessary



# $1\pi$ production

## ■ Reactions

■ Incoherent:  $\nu_l A \rightarrow l \pi X$

■ Coherent:

■ CC  $\nu_l A \rightarrow l^- \pi^+ A$

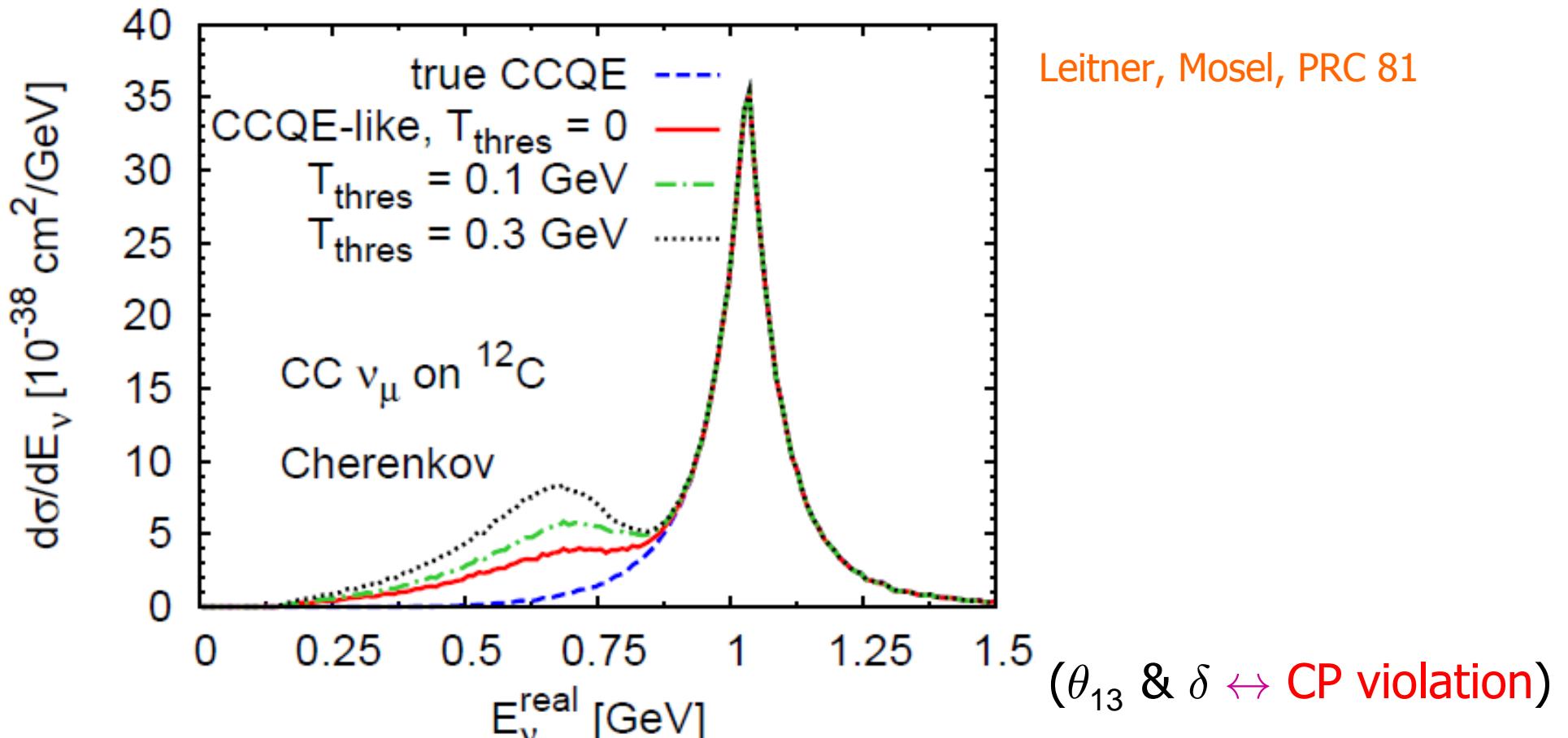
■ NC  $\nu A \rightarrow \nu \pi^0 A$

## ■ Relevant for oscillations

■ NC  $\pi^0$ : e-like background to  $\nu_\mu \rightarrow \nu_e$  searches ( $\theta_{13}$  &  $\delta \leftrightarrow$  CP violation)

■ Source of CCQE-like events (in nuclei), needs to be subtracted for a good  $E_\nu$  reconstruction

# $1\pi$ production



- Source of CCQE-like events (in nuclei), needs to be subtracted for a good  $E_\nu$  reconstruction

# 1 $\pi$ production

- Elementary process:  $\nu_l \ N \rightarrow l \ \pi \ N'$

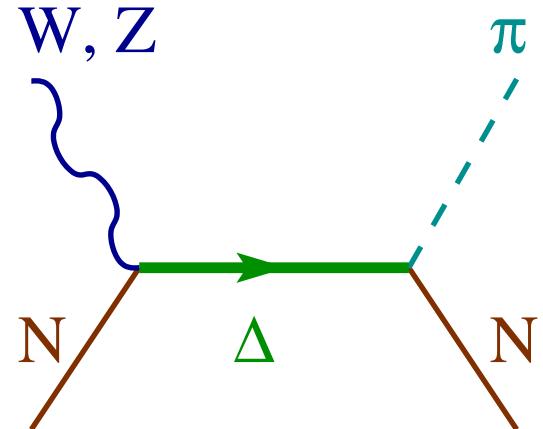
- Dominated by resonance production

- At  $E_\nu \sim 1$  GeV:  $\Delta(1232)$

$$\mathcal{M} = \frac{G_F \cos \theta_C}{\sqrt{2}} l^\alpha J_\alpha$$

- N- $\Delta$  transition current:

$$J^\mu = \bar{\psi}_\mu \left[ \left( \frac{C_3^V}{M} (g^{\beta\mu} q - q^\beta \gamma^\mu) + \frac{C_4^V}{M^2} (g^{\beta\mu} q \cdot p' - q^\beta p'^\mu) + \frac{C_5^V}{M^2} (g^{\beta\mu} q \cdot p - q^\beta p^\mu) \right) \gamma_5 + \frac{C_3^A}{M} (g^{\beta\mu} q - q^\beta \gamma^\mu) + \frac{C_4^A}{M^2} (g^{\beta\mu} q \cdot p' - q^\beta p'^\mu) + C_5^A g^{\beta\mu} + \frac{C_6^A}{M^2} q^\beta q^\mu \right] u$$



- Form factors  $\Leftrightarrow$  Helicity amplitudes ( $A_{1/2}$ ,  $A_{3/2}$ ,  $S_{1/2}$ )

# $1\pi$ production

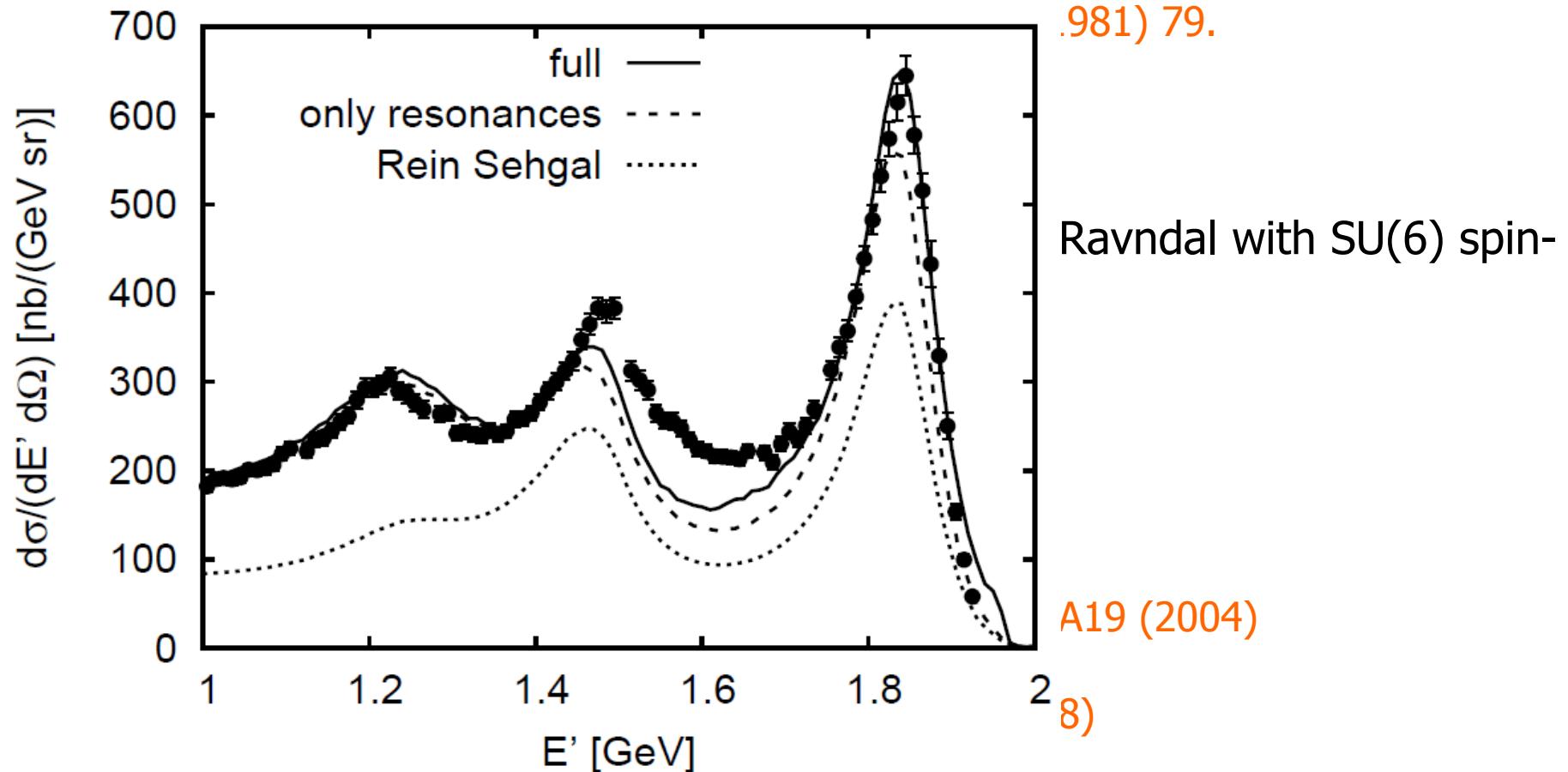
- Elementary process:  $\nu_l N \rightarrow l \pi N'$ 
  - Dominated by resonance production
  - Rein-Sehgal model: Rein-Sehgal, Ann. Phys. 133 (1981) 79.
    - Used by almost all MC generators
    - Relativistic quark model of Feynman-Kislinger-Ravndal with SU(6) spin-flavor symmetry
    - Helicity amplitudes for 18 baryon resonances
    - Lepton mass = 0
      - Corrections: Kuzmin et al., Mod. Phys. Lett. A19 (2004)  
Berger, Sehgal, PRD 76 (2007)  
Graczyk, Sobczyk, PRD 77 (2008)
    - Poor description of  $\pi$  electroproduction data on p

# $1\pi$ production

- Elementary process:  $\nu_l N \rightarrow l \pi N'$

- Dominated by resonance production

$e^- p \rightarrow e^- X, \theta = 20^\circ, E=2.445 \text{ GeV}$



- Poor description of  $\pi$  electroproduction data on p

# $1\pi$ production

## ■ $N-\Delta$ transition current

$$J^\mu = \bar{\psi}_\mu \left[ \left( \frac{C_3^V}{M} (g^{\beta\mu} q^\ell - q^\beta \gamma^\mu) + \frac{C_4^V}{M^2} (g^{\beta\mu} q \cdot p' - q^\beta p'^\mu) + \frac{C_5^V}{M^2} (g^{\beta\mu} q \cdot p - q^\beta p^\mu) \right) \gamma_5 \right. \\ \left. + \frac{C_3^A}{M} (g^{\beta\mu} q^\ell - q^\beta \gamma^\mu) + \frac{C_4^A}{M^2} (g^{\beta\mu} q \cdot p' - q^\beta p'^\mu) + C_5^A g^{\beta\mu} + \frac{C_6^A}{M^2} q^\beta q^\mu \right] u$$

## ■ Helicity amplitudes can be extracted from data on $\pi$ photo- and electro-production

- Unitary isobar model MAID Drechsel, Kamalov, Tiator, EPJA 34 (2007) 69
  - Uses world data
  - for all 4 star resonances with  $W < 2$  GeV
- Unitary isobar model+Regge-pole BG at high energies I. Aznauryan, PRC67
- Dispersion relations
  - CLAS (JLab) data
  - 1<sup>st</sup> and 2<sup>nd</sup> resonance regions:  $\Delta(1232)$ ,  $N^*(1440)$ ,  $N^*(1520)$ ,  $N^*(1535)$

# $1\pi$ production

## ■ $N-\Delta$ transition current

$$J^\mu = \bar{\psi}_\mu \left[ \left( \frac{C_3^V}{M} (g^{\beta\mu} q^\not - q^\beta \gamma^\mu) + \frac{C_4^V}{M^2} (g^{\beta\mu} q \cdot p' - q^\beta p'^\mu) + \frac{C_5^V}{M^2} (g^{\beta\mu} q \cdot p - q^\beta p^\mu) \right) \gamma_5 \right. \\ \left. + \frac{C_3^A}{M} (g^{\beta\mu} q^\not - q^\beta \gamma^\mu) + \frac{C_4^A}{M^2} (g^{\beta\mu} q \cdot p' - q^\beta p'^\mu) + C_5^A g^{\beta\mu} + \frac{C_6^A}{M^2} q^\beta q^\mu \right] u$$

## ■ Axial form factors

$$C_6^A = C_5^A \frac{M^2}{m_\pi^2 + Q^2} \leftarrow \text{PCAC}$$

$$C_4^A = -\frac{1}{4} C_5^A \quad C_3^A = 0 \leftarrow \text{Adler model}$$

$$C_5^A = C_5^A(0) \left( 1 + \frac{Q^2}{M_{A\Delta}^2} \right)^{-1}$$

# $1\pi$ production

## ■ $N-\Delta$ axial form factors: determination of $C_A^5(0)$ and $M_{A\Delta}$

■  $C_5^A(0) = \frac{g_\Delta N_\pi f_\pi}{\sqrt{6} M} \approx 1.2 \leftarrow$  off diagonal GT relation

■ From ANL and BNL data on  $\nu_\mu d \rightarrow \mu^- \pi^+ p n$

■ with large normalization (flux) uncertainties

■ Graczyk et al., PRD 80 (2009)

■ Deuteron effects

■ Non-resonant background absent

■  $C_A^5(0) = 1.19 \pm 0.08$ ,  $M_{A\Delta} = 0.94 \pm 0.03$  GeV

■ Hernandez et al., PRD 81 (2010)

■ Deuteron effects

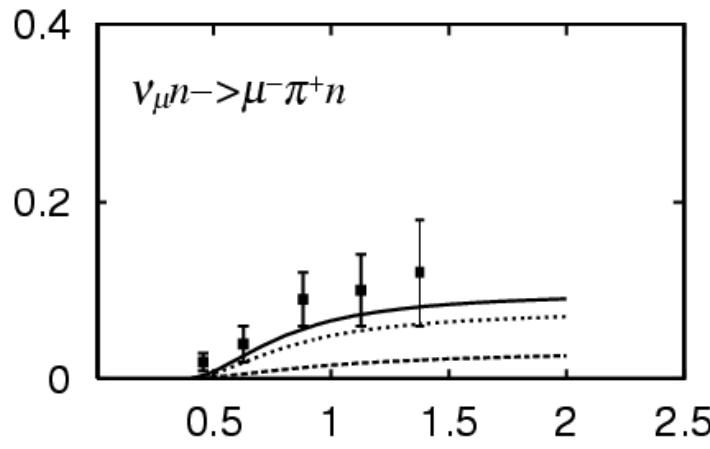
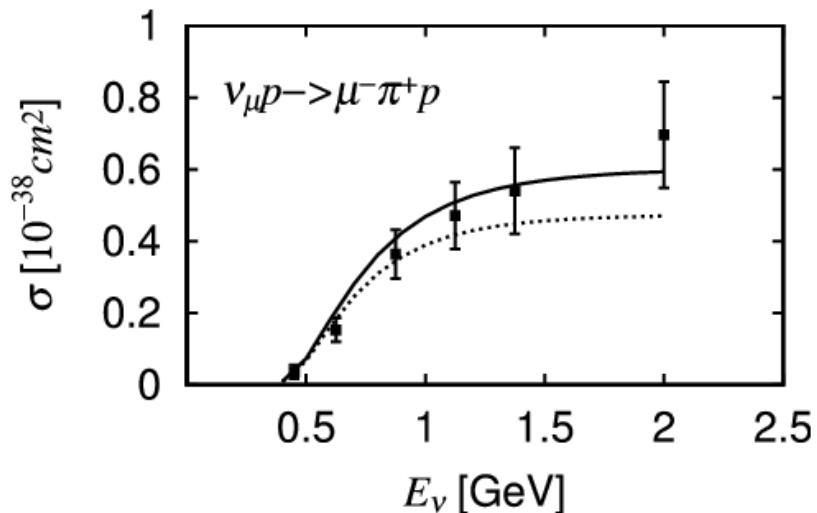
■ Non-resonant background fixed by chiral symmetry

■  $C_A^5(0) = 1.00 \pm 0.11$  GeV,  $M_{A\Delta} = 0.93 \pm 0.07$  GeV

■ 20 % reduction of the GT relation

# $1\pi$ production

- Elementary process
  - Sato & Lee model [Sato, Uno, Lee, PRC 67 \(2003\)](#)
  - Dynamical model for  $\pi$  production with  $\gamma, e, \nu$
  - Starting with an effective H:  $\pi N, \Delta N \Rightarrow$
  - T-matrix obtained from coupled channel [Lippman-Schwinger](#) eq.
  - Good agreement with data



- Bare  $\Delta N$  renormalized by meson clouds (30 %):  $C_{A_5}^{A_5}(0) = 0.96 \text{ GeV}$  reconciles the empirical value with [quark model](#) results

# $1\pi$ production

- Incoherent  $1\pi$  production in nuclei
  - Large number of excited states  $\Rightarrow$  semiclassical treatment
  - $\pi$  propagation (scattering, charge exchange), absorption (FSI)
  - Most models cannot calculate this reaction channel.

Exceptions:

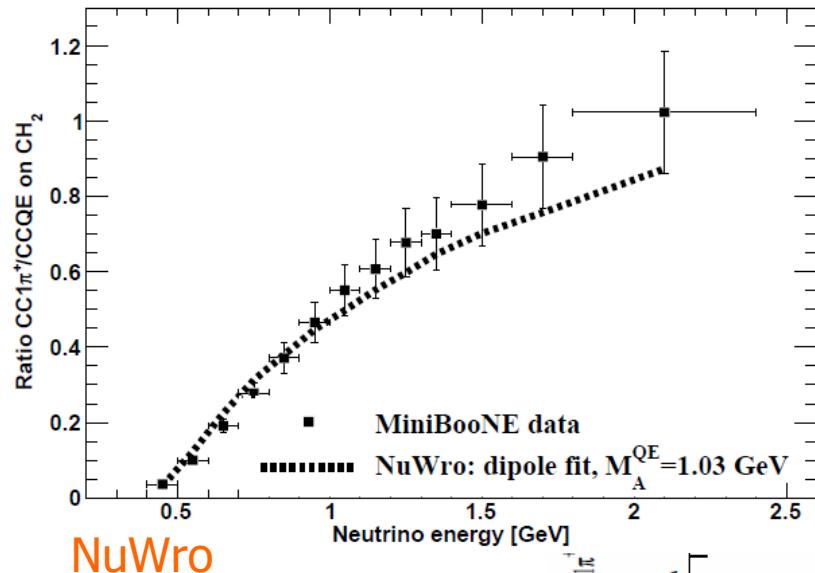
- MC generators: NUANCE, NEUT, GENIE, NuWro
- Cascade: Ahmad et al., PRD 74 (2006)
- Transport: GiBUU

# $1\pi$ production

- Incoherent  $1\pi$  production in nuclei (FSI)
  - NuWro J. Sobczyk et al.
    - Intranuclear cascade
    - $\pi$  propagation: empirical  $\pi$ -N vacuum  $\sigma$
    - $\pi$  absorption:  $\pi$ -A absorption data
  - Ahmad et al., PRD 74 (2006)
    - Cascade ( $\sim$ NuWro)
    - In-medium modification of  $\Delta$  spectral f. (only in the production)
  - GiBUU
    - Transport: one approach for eA,  $\nu$ A, pA,  $\pi$ A reactions
    - $\pi$ , N but also  $\Delta$  are propagated
    - Main absorption mech.:  $\Delta$  N  $\rightarrow$  N N,  $\pi$  N N  $\rightarrow$  N N

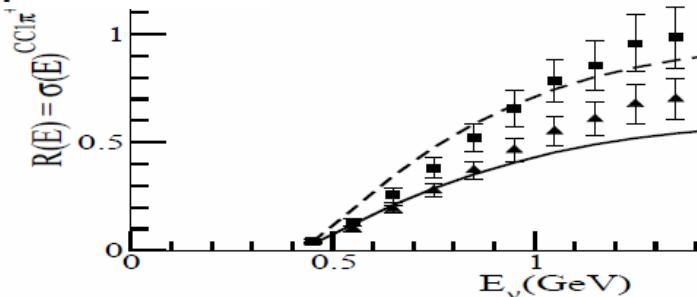
# $1\pi$ production

- Comparison to the  $\sigma(\text{CC}\pi^+)/\sigma(\text{CCQE-like})$  ratio at MiniBooNE

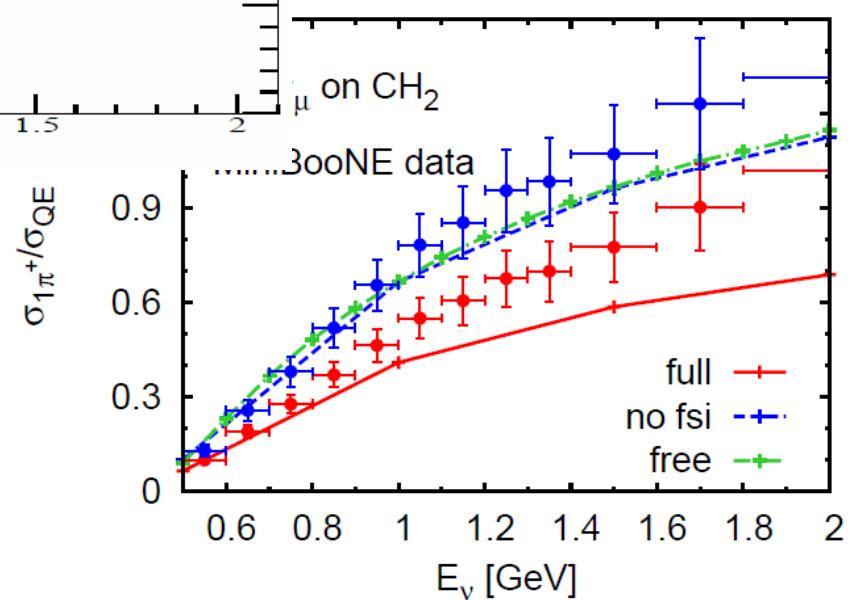


Athar et al.

1pi+(w/o ME) / CCQE (w/o RPA)  
1pi+ (with ME + FSI effects) / CCQE (with RPA+ QE-like)  
1pi+/CCQE (FSI Corrected)  
1pi+-like/CCQE-like(observed)



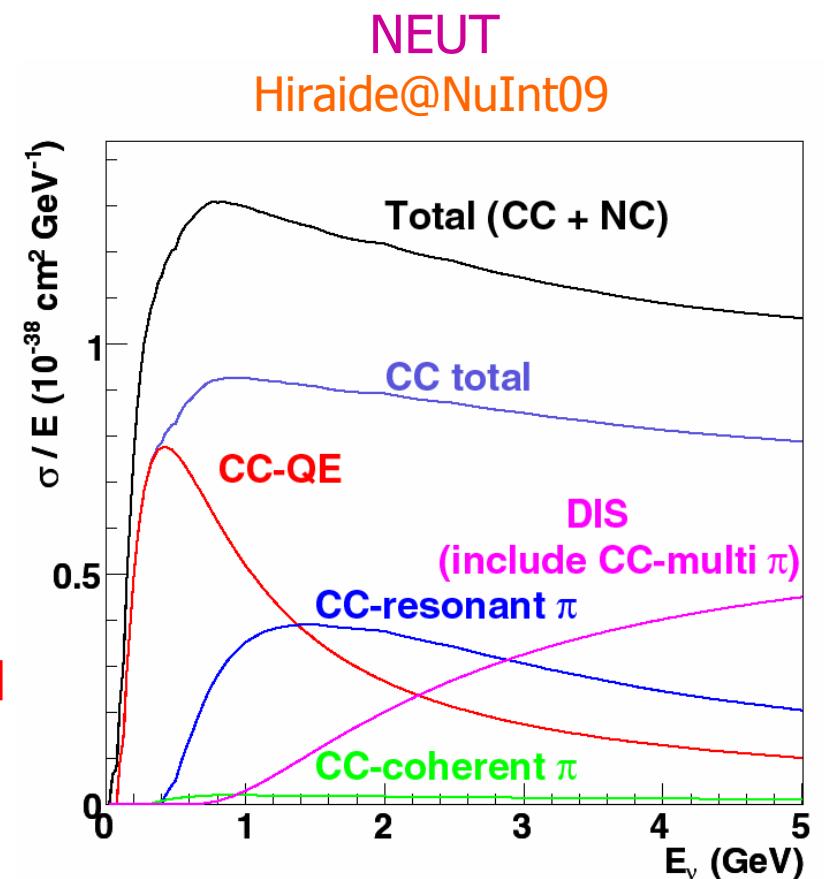
GiBUU



# $1\pi$ production

## ■ Coherent pion production

- CC  $\nu_l A \rightarrow l^- \pi^+ A$
- NC  $\nu A \rightarrow \nu \pi^0 A$
- Takes place at low  $q^2$
- Very small cross section
- At  $q^2 \sim 0$ , axial current **not suppressed**



# $1\pi$ production

## ■ Coherent pion production models

### ■ PCAC Rein & Seagal NPB 223 (83) 29

Berger & Seagal, PRD 76 (2007), 79 (2009)

Paschos & Schalla, PRD 80 (2009),

- In the  $q^2=0$  limit, PCAC is used to relate  $\nu$  induced coherent pion production to  $\pi A$  elastic scattering

- Extrapolated to  $q^2 \neq 0$

- R&S: describe  $\pi A$  in terms of  $\pi N$  scattering

- B&S, P&S: use  $\pi A$  data

### ■ Problems: Hernandez et al., PRD 80 (2009) 013003

- $q^2=0$  limit neglects important angular dependence at low energies

- R&S: The  $\pi A$  elastic description is not realistic

- B&S,P&S: spurious initial  $\pi$  distortion present in  $\pi A$  but not in coh $\pi$

# $1\pi$ production

## ■ Coherent pion production models

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- R&S: describe  $\pi A$  in terms of  $\pi N$  scattering

- B&S, P&S: use  $\pi A$  data

- Problems of PCAC models: less relevant as the energy increases

- NOMAD:  $\sigma = 72.6 \pm 8.1(\text{stat}) \pm 6.9(\text{syst}) \times 10^{-40} \text{ cm}^2$

- Consistent with RS:  $\sigma \approx 78 \times 10^{-40} \text{ cm}^2$

# $1\pi$ production

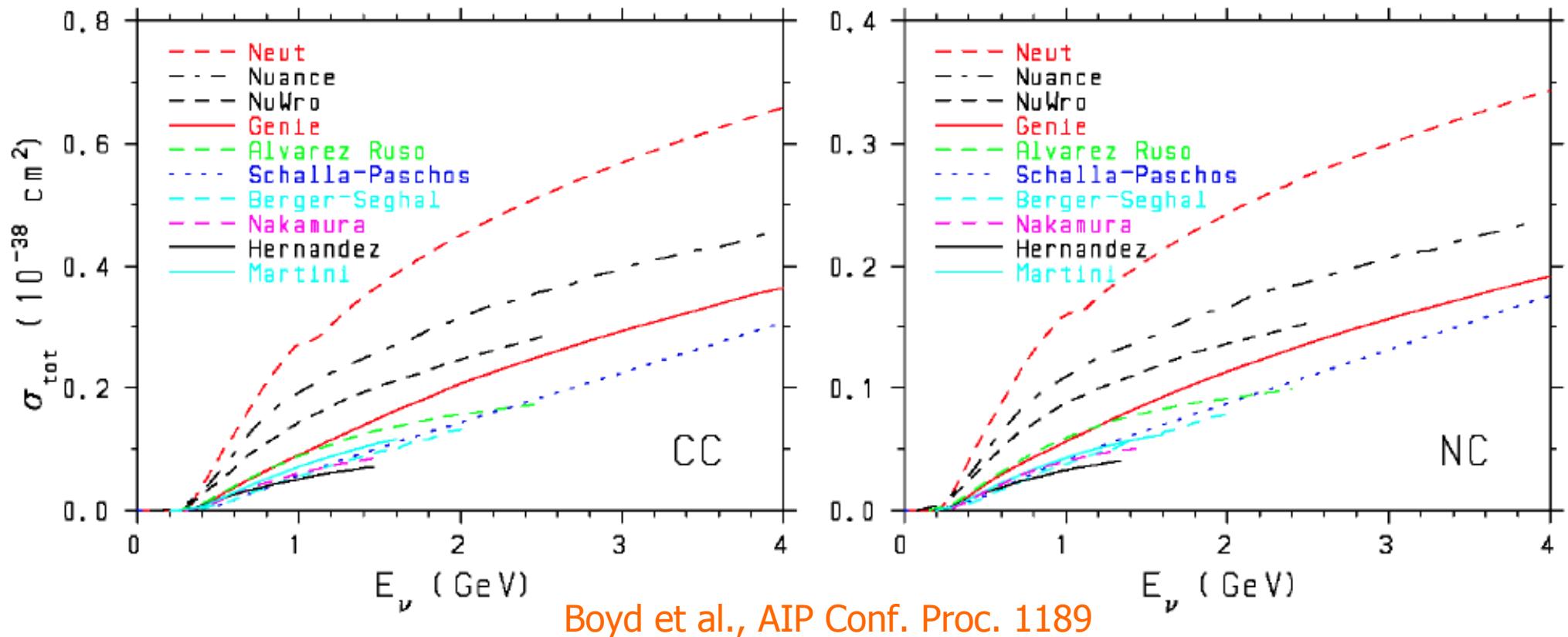
- Coherent pion production

- Microscopic models: Singh et al., PRL 96 (2006)  
LAR et al, PRC 76 (2007)  
Amaro et al., PRD 79 (2009)  
Nakamura et al., PRC 81 (2010)

- $\Delta$  excitation is dominant
  - $\Delta$  properties change in the nuclear medium
  - $\pi$  distortion: DWIA with optical potential based on  $\Delta$ -hole model
  - Treatment is consistent with incoherent  $\pi$  production
  - Valid only at low energies
  - $\sigma \sim [C_5^A(0)]^2$

# $1\pi$ production

## ■ Coherent pion production

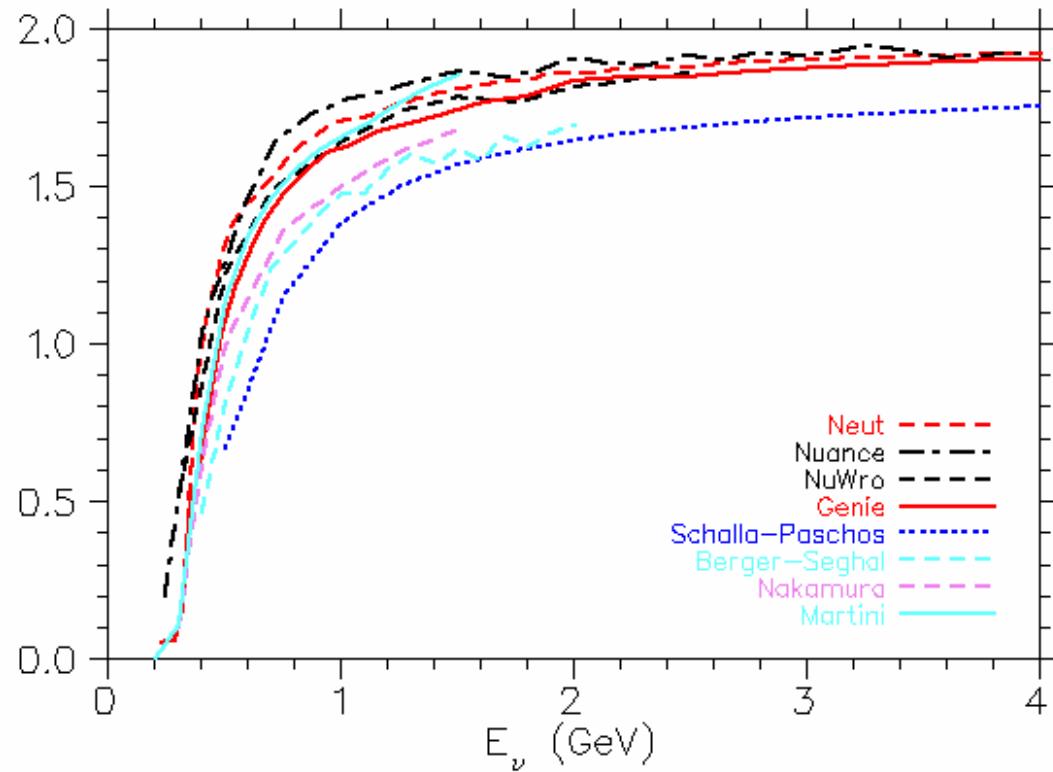


# $1\pi$ production

## ■ Coherent pion production

Ratio of CC to NC total

Boyd et al., AIP Conf. Proc. 1189



- SciBooNE: PRD 81 (2010)
- NC  $\pi^0 \sigma$  compatible with R&S
- $CC\pi^+/NC\pi^0 = 0.14^{+0.30}_{-0.28}$
- Theoretical models predict  $CC\pi^+/NC\pi^0 \sim 1-2$  !

# Conclusions

- In spite of the exp. and th. effort **basic  $\nu A$**  interaction mechanisms (QE,  $1\pi$ ) are not understood
- Comparison to **inclusive** data is needed
- Look forward to **new** data and theoretical progress